

Lecture 26
Turbulence - Chemistry

"Prepols" not working

Chemistry-Turbulence Interactions

1. Idea of time scales (from Smeets & Smith, p. 269-273)
2. Magnussen-Hjertager approach from {Fluent Manual,
Magnussen & Hjertager, 16th Symp., 1976
Smith & Fletcher, CST, 1988}
3. Mixture fraction approach
4. Assumed-shape PDF
5. Smith & Fletcher paper

THE PROBLEM

Species continuity eqn:

$$\frac{d}{dt} (y_i \rho) + \nabla \cdot (\rho u y_i - \bar{D}_{y_i} \nabla y_i) = R_{y_i}$$

Introduce time-mean (Reynolds) averaging

$$y_i = \bar{y}_i + y_i' \quad \bar{y}_i = \frac{1}{\Delta t} \int_0^{\Delta t} y_i dt \quad (\text{short time})$$

For now, assume density fluctuations are small, $\overline{y_i' u'} \approx 0$

$$\frac{d}{dt} (\bar{\rho} \bar{y}_i) + \nabla \cdot (\bar{\rho} \bar{u} \bar{y}_i - \bar{D}_{y_i} \nabla \bar{y}_i) = \bar{R}_{y_i}$$

HOWEVER) $\bar{R}_{y_i} \neq R_{\bar{y}_i}$

$$\begin{aligned} \bar{R}_{y_i} &= \frac{1}{\Delta t} \int_0^{\Delta t} R_{y_i} dt \\ &= \frac{1}{\Delta t} \int_0^{\Delta t} \sum_{\text{rxns}} \nu_i A_i e^{-E_i/RT} y_{i1} y_{i2} dt \end{aligned}$$

① Very sensitive to fluctuations in T② cannot ignore $\overline{y_i' T'}$ correlations③ bad to assume $\bar{R}_{y_i} = \sum \nu_i A_i e^{-E_i/\bar{R}\bar{T}} \bar{y}_{i1} \bar{y}_{i2}$

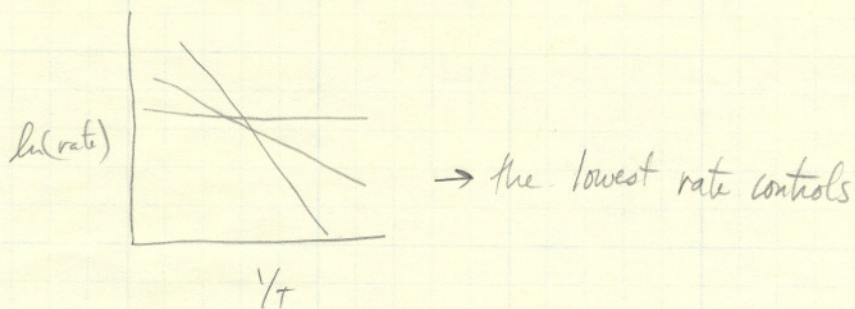
(although some people have done this)

④ fluctuations in T affect energy equation as well

 $\Rightarrow$ Need for some slick model!!

Concept of time scales

A. 1st think of reaction rates (like of a solid)



B. Rate of mixing vs. rate of reaction → lowest rate controls

τ_r = reaction time scale (time to react ~ 90%)

τ_m = mixing time scale (time to mix on a molecular level)

Case 1

$$\tau_m \ll \tau_r$$

$$\bar{R}_i = f(\bar{C}_i, \bar{C}_j, \bar{p}, \bar{T}) = A e^{-\frac{E}{RT}} \bar{C}_i^\alpha \bar{C}_j^\beta$$

Case 2

$$\tau_m \gg \tau_r$$

- Calculate mixing rate only
- If its mixed, its burned
- equilibrium kinetics after mixing
- not valid for premixed systems

Case 3

$$\tau_m \approx \tau_r$$

- Difficult to handle

What usually happens in combustion?

Case 3 (or case 2 for diffusion flames)

Magnussen-Hjertager 16th Symp., 1976

Proposal: Calculate (a) the mixing rate
(b) the reaction rate based on mean values

Take the lower of (a) & (b)

Fuel + Oxidizer $\rightarrow$ Products

A. Mixing-limited rate of fuel consumption:

$$(1) \quad R_f = A \bar{C}_f \left(\frac{\bar{E}}{k} \right)$$

$\uparrow$ empirical constant $\uparrow$ mixing rate

B. Mixing-limited rate of oxidizer consumption

$$(2) \quad R_f = A \frac{\bar{C}_{O_2}}{r_f} \left(\frac{\bar{E}}{k} \right), \quad r_f = \text{stoichiometric oxygen requirement}$$

C. For premixed flames, you need the rate of mixing of hot products

$$(3) \quad R_f = A \cdot B \cdot \frac{\bar{C}_p}{1+r_f} \frac{\bar{E}}{k}$$

D. Arrhenius reaction rate

$$(4) \quad R_f = k_0 e^{-E/(RT)} C_f^a C_o^b$$

$\Rightarrow$ Actual Rate = minimum of {1, 2, 3, 4}

$\rightarrow$ You either get reactants or products (no intermediates, like CO)

$\rightarrow$ one reaction may not be totally controlling

$\rightarrow$ multiple reactions? not!

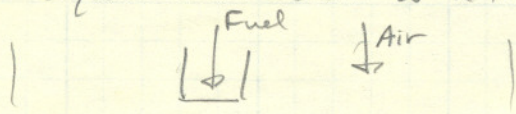
$\rightarrow$ fast! (computationally)

Mixing-limited reactions

People have found experimentally that species concentrations and temperatures in laminar diffusion flames are functions of the local equivalence ratio.

✶ - Picture from Smart & Smith

Equivalence Ratio is a function of Mixture Fraction



$$f = \frac{m_p}{m_p + m_s}$$

Any conserved scalar "s" can then be calculated

$$S = f s_p + (1-f) s_s$$

↑
value of S in
pure
primary

Examples of conserved scalar :- mass of element i (1-phase)

- density (non-reacting only)
- enthalpy (adiabatic only)

Idea: Calculate f , then get all of the species!!

Conservation equation for " $\tilde{f}$ ":

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{f}) + \vec{\nabla} \cdot (\bar{\rho} \tilde{f} \vec{v} - \mathcal{D}_{\text{turb}, f} \vec{\nabla} \tilde{f}) = 0$$

$$D_{\text{turb}, p} = \frac{v_g^t}{\sigma_f}$$

Smoot; Smith
11.21

Also, define $\hat{g}_f = \frac{(p_f - \bar{p}_f)^2}{\bar{p}}$ (density weighted)

Transport eqn:

$$\frac{\partial}{\partial t}(\rho T_f) + \nabla \cdot (\rho \vec{v} T_f - \frac{\rho}{g} \nabla T_f) = C_{f1} \mu_t \left[\left(\frac{\partial \tilde{f}}{\partial x} \right)^2 + \left(\frac{\partial \tilde{f}}{\partial r} \right)^2 \right] - \frac{C_{f2} \tilde{\rho} \tilde{\epsilon} T_f}{k^2}$$

generation dissipation

Sweet
smith

11.29

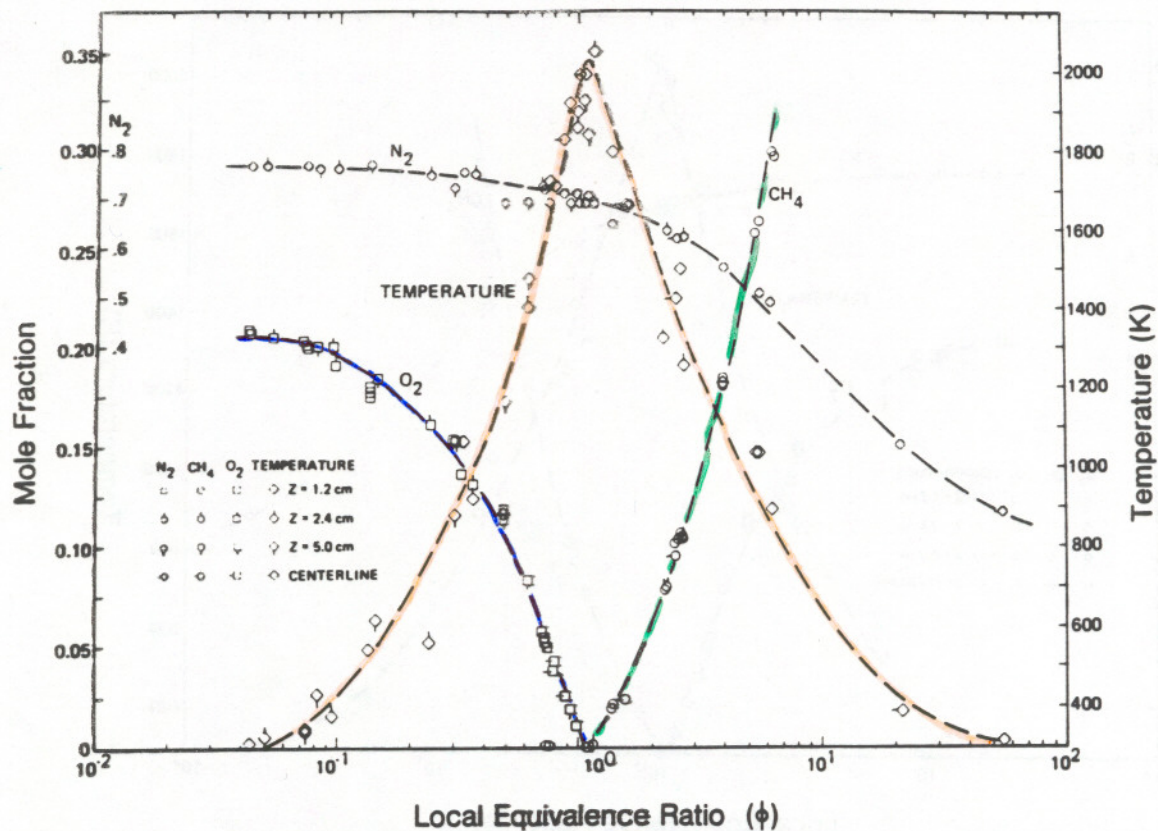


Figure 11.4. Measured temperature and reactant concentration as a function of the local equivalence ratio in a laminar methane diffusion flame. (Figure used with permission from Mitchell *et al.*, 1980.)

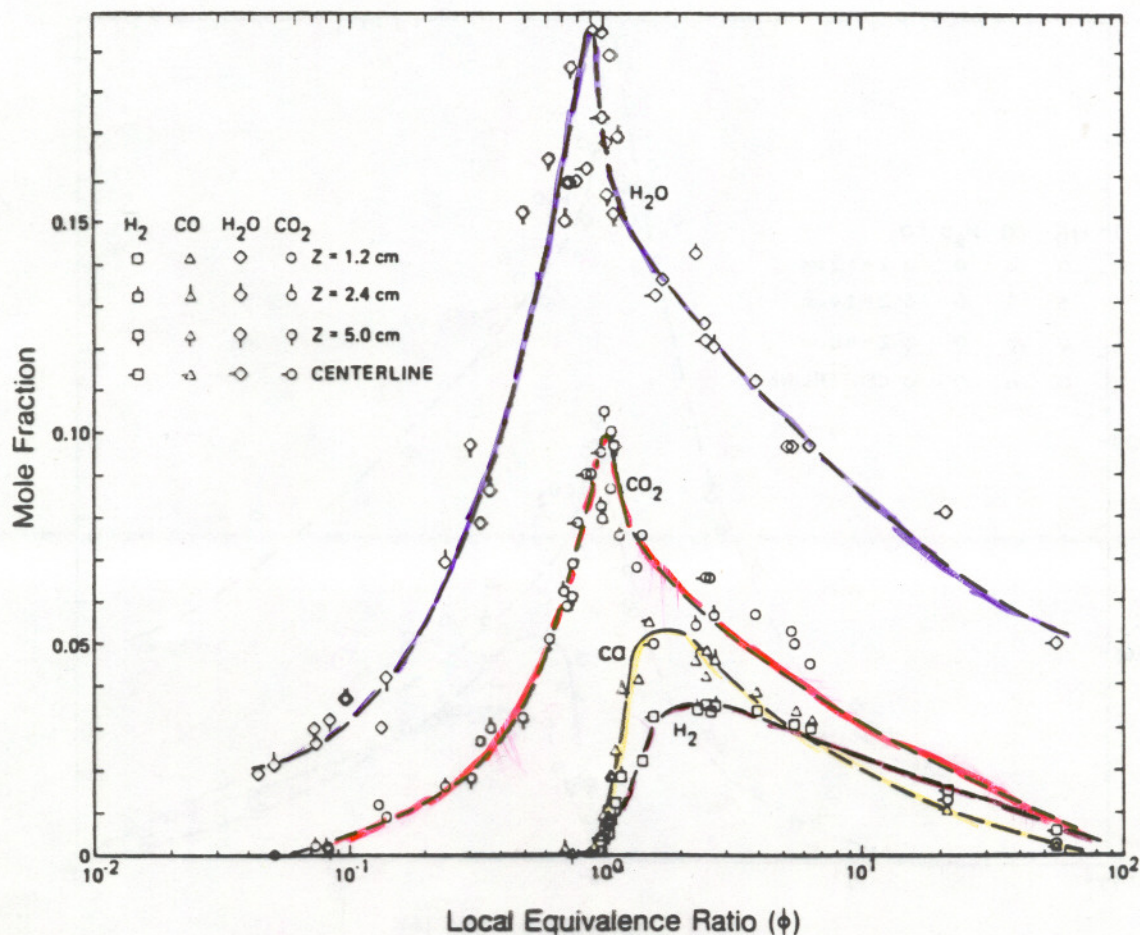


Figure 11.5. Measured major-product concentration as a function of the local-equivalence ratio in a laminar methane diffusion flame. (Figure used with permission from Mitchell *et al.*, 1980.)

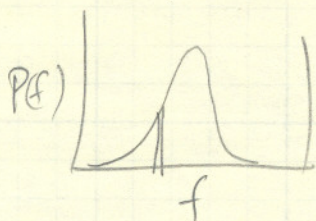
Lecture 27

PDF's for

turbulent combustion

(PDF's)

$P(f) df = \text{fraction of time you are at } f$



Properties of PDF's (Identities)

$$\int_0^1 P(f) df = 1$$

$$\bar{f} = \int_0^1 f P(f) df$$

$$g_f = \int_0^1 (f - \bar{f})^2 P(f) df$$

Idea: ① Assume gaussian shape, which

is defined by $\bar{f}$ & g_f

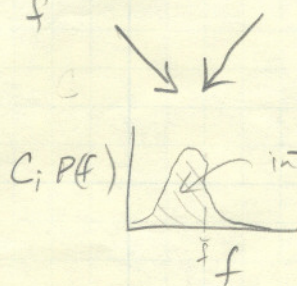
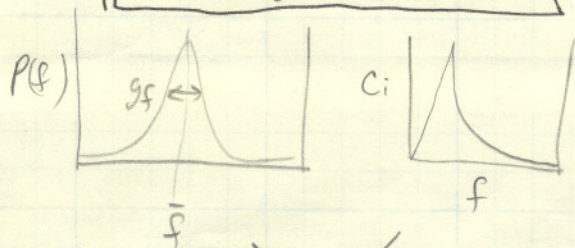
$$P(f) = \frac{1}{\sqrt{2\pi}g_f} e^{-\frac{(f-\bar{f})^2}{2g_f}}$$

② Instantaneous properties (species) can be obtained for each instantaneous f [table of y_i vs f]

③ Mean properties can be obtained using the PDF (use equilibrium based on elemental comp. and n)

$$\bar{\beta} = \int_0^1 \beta P(f) df$$

if $\beta = f(\text{elemental comp., enthalpy})$
 $= f(f)$



integral =

$$\bar{C}_i$$

wanted!

- (more important than rxn rate!)
- (species const. eqns. not solved)

PDFAdvantages

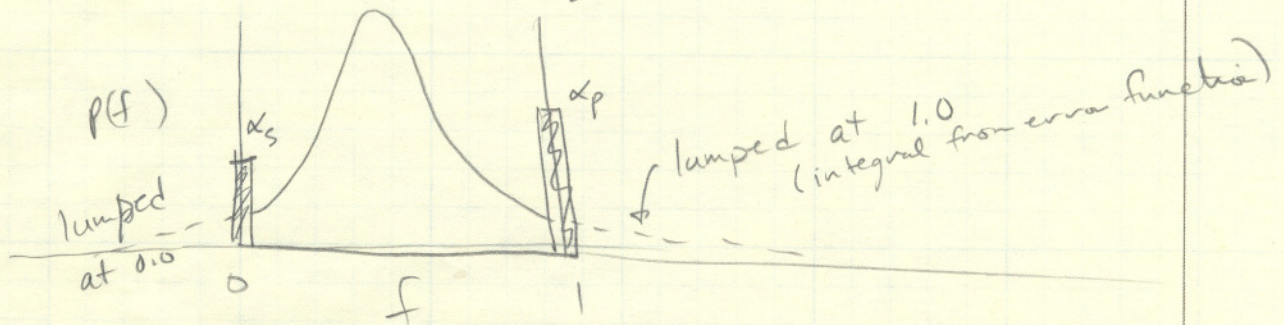
- no reaction rates required (use equilibrium & data)
- intermediate species predicted
- only one transport equation!

Disadvantages

- non-premixed flames only
- reactions may not be mixing-limited in some cases, which goes against assumptions

Other: you may have pockets of pure, unreacted primary or secondary

⇒ intermittency



readjust integrals so

$$\alpha_p + \alpha_s + \int_0^1 P(f) df = 1$$

$$\alpha_p f_p + \alpha_s f_s + \int_0^1 f P(f) df = \bar{f}$$

also $\int f df \Rightarrow$ ^{must be} correct (remember $f_s = 0$, $f_p = 1$)

★ ⇒ Table 11.2 in Smoot & Smith

TABLE 11.2. Parameters for Clipped Gaussian Probability Density Function

$$P(f) = (2\pi G_f)^{-1/2} \exp(-Z_f^2/2)$$

$$\alpha = (2\pi)^{-1/2} \int_L^U \exp(-Z_f^2/2) dZ_f$$

Intermittency	U	L	Z_f
α_p	∞	$(1-F)/G_f^{1/2}$	$(f-F)/G_f^{1/2}$
α_s	$-F/G_f^{1/2}$	$-\infty$	$(f-F)/G_f^{1/2}$

where F and G_f come from

$$\bar{f} = \alpha_p + (2\pi G_f)^{-1/2} \int_{0+}^{1-} f \exp\left(\frac{-(f-F)^2}{2G_f}\right) df$$

$$g_f = \alpha_p - \bar{f}^2 + (2\pi G_f)^{-1/2} \int_{0+}^{1-} f^2 \exp\left(\frac{-(f-F)^2}{2G_f}\right) df$$

Example

CH₄/air at 300 K

Primary = CH₄ @ 300 K, $H_p = -17.889 \frac{\text{kcal}}{\text{gmole}}$

Secondary = Air @ 300 K, $H_s = 0$



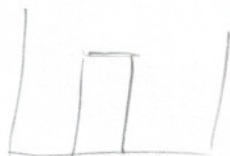
Say we have 1 kmole/s of CH₄
Need 2 kmole/s of O₂

⇒ See computer disk

Fluent ⇒ uses a B-function instead of a Gaussian
(no exponential involved)

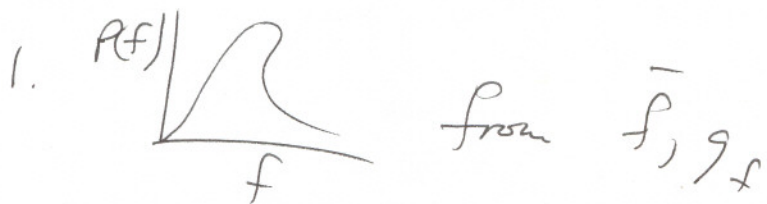
$$\text{pdf} = \frac{f^a (1-f)^b}{B}$$

⇒ Top hat can be used



⇒ stores completely integrated PDF stuff in a table

f	$g(f)$	$\bar{T}$	$\bar{Y}_{\text{O}_2}$	etc.
$\vdots$	$\vdots$	$\vdots$	$\vdots$	



2. y_k at different values of f

f	y_k 's	h	T_{eq}	y_i
			—	—
			—	—
			—	—
			—	—

3. $\int T(f) P(f) df$

T

