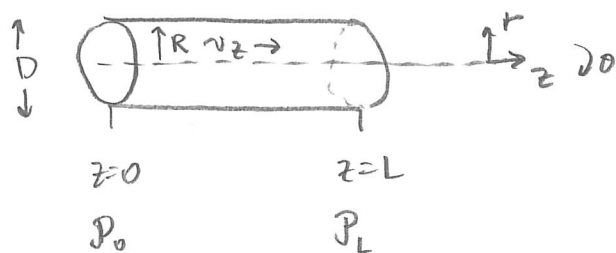


Velocity Profile near a wall in turbulent flow



$$\nabla \cdot \langle \underline{u} \rangle = 0, \quad \nabla \cdot \underline{u} = 0$$

* unidirectional for
 $\langle \underline{u} \rangle$ in $\underline{z}$
 (b/c of average)

$$\rho \frac{D\langle \underline{u} \rangle}{Dt} = - \nabla \langle P \rangle + \nabla \cdot (\langle \underline{\tau} \rangle + \underline{\tau}^*)$$

For the given geometry, this simplifies to:

$$0 = - \frac{\partial \langle P \rangle}{\partial z} + \frac{1}{r} \frac{\partial}{\partial r} \left[r (\langle \tau_{rz} \rangle + \tau_{rz}^*) \right]$$

recall that: $\tau_w = -\frac{D}{4} \frac{\Delta P}{L} = -\frac{D}{4} \frac{\partial \langle P \rangle}{\partial z}$ (Eq. 2.3-4 on p. 35)

$$\frac{\partial \langle P \rangle}{\partial z} = -\frac{4\tau_w}{D} = -\frac{2}{R} \tau_w$$

$R = D/2$

↑ a constant.

$$\frac{2}{R} \tau_w = \frac{1}{r} \frac{\partial}{\partial r} \left[r (\langle \tau_{rz} \rangle + \tau_{rz}^*) \right]$$

↓
 let $\tau_0 = -\tau_w$ because
 this is what the book defines.

$$\frac{2r}{R} \tau_0 = \frac{\partial}{\partial r} \left[r (\langle \tau_{rz} \rangle + \tau_{rz}^*) \right]$$

↓ separate & integrate

$$\int \frac{2r}{R} \tau_0 dr = \int d \left[r (\langle \tau_{rz} \rangle + \tau_{rz}^*) \right]$$

$$\frac{2r^2}{2R} T_0 + C = r (\langle \tau_{rz} \rangle + \tau_{rz}^*)$$

$$\frac{r}{R} T_0 + \frac{C}{r} = \langle \tau_{rz} \rangle + \tau_{rz}^*$$

via symmetry

$$C=0 \begin{cases} \langle \tau_{rz} \rangle(0) + \tau_{rz}^*(0) = 0 \\ \text{or} \\ \langle \tau_{rz} \rangle(R) + \tau_{rz}^*(R) = T_0 \end{cases}$$

$$\langle \tau_{rz} \rangle + \tau_{rz}^* = \frac{T_0 r}{R}$$

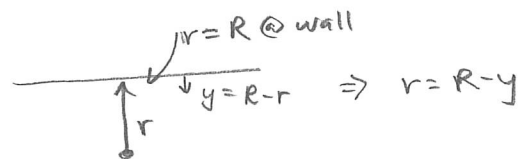
← (Eq. 10.5-4), but book has a sign error.

* change of variables

$$\text{let } y = R - r$$

(the part near the wall)

and assume $y \ll R$.



$$\langle \tau_{yz} \rangle + \tau_{yz}^* = \frac{T_0}{R} (R - y)$$

$$= T_0 (1 - y/R)$$

* change to look @ cartesian

T only. This is okay because

$y \ll R$, so the radius of

curvature is negligible. →

now we can say

$$\langle \tau_{yz} \rangle = \rho \nu \frac{d\langle v_z \rangle}{dy}$$

$$\tau_{yz}^* = \rho \epsilon \frac{d\langle v_z \rangle}{dy}$$

↑
eddy diffusivity

$$\rho(y+\varepsilon) \frac{\langle v_z \rangle}{dy} = \tau_0 \quad \leftarrow \text{the book makes this dimensionless.}$$

$$\text{BC: } \langle v_z \rangle(0) = 0 \quad (\text{no slip @ the wall})$$

$$v_z^+ = \frac{\langle v_z \rangle}{(\tau_0/\rho)^{1/2}} = \frac{\langle v_z \rangle}{u_\tau}$$

$$\left(\frac{\text{kg m}}{\text{s}^2 \text{ m}^2} \cdot \frac{1}{\text{m}^2} \right)^{1/2} = \left(\frac{\text{m}^3}{\text{s}^2 \text{ kg}} \right)^{1/2} = \left(\frac{\text{m}^2}{\text{s}^2} \right)^{1/2} = \frac{\text{m}}{\text{s}}$$

v_z^+, y^+
are called
"wall variables"

$$u_\tau = (\tau_0/\rho)^{1/2} = \text{"friction velocity"}$$

$$y^+ = y / (v/u_\tau)$$

$$v/u_\tau = \frac{\text{m}^2}{\text{s}} \cdot \frac{\text{s}}{\text{m}} = \text{m}$$

$$\rho(y+\varepsilon) \frac{u_\tau}{v} u_\tau \frac{dv_z^+}{dy^+} = \tau_0$$

$$\frac{u_\tau^2 \rho}{v} (y+\varepsilon) \cdot \frac{1}{\tau_0} \frac{dv_z^+}{dy^+} = 1 \quad \leftarrow u_\tau^2 = \frac{\tau_0}{\rho}$$

$$\frac{\tau_0}{\rho} \frac{\rho}{v} \cdot \frac{1}{\tau_0} (y+\varepsilon) \frac{dv_z^+}{dy^+} = 1$$

$$\boxed{\left(1 + \frac{\varepsilon}{y}\right) \frac{dv_z^+}{dy^+} = 1, \quad v_z^+(0) = 0}$$

(Eq. 10.5-7)

Now it is math
from here.

This is not as easy as it looks because
the eddy diffusivity is not a constant

$$\frac{\varepsilon}{v} = (k y^+)^2 \left(\frac{dv_z^+}{dy^+} \right) \quad \leftarrow \text{Eq. 10.4-9.}$$

↑
constant that is empirically fit to data.

* when $y^+ \rightarrow 0$, $\frac{\varepsilon}{\nu} \rightarrow 0$ (inner region)

$$\frac{dv_z^+}{dy^+} = 1 \quad v_z^+(0) = 0 \Rightarrow \boxed{v_z^+ = y^+ \quad (y^+ \rightarrow 0)}$$

* when $y^+ \rightarrow \infty$, $(ky^+)^2 \frac{dv_z^+}{dy^+} \gg 1$ (outer region)

$$\left(ky^+ \frac{dv_z^+}{dy^+} \right)^2 = 1$$

$$ky^+ \frac{dv_z^+}{dy^+} = 1 \Rightarrow \frac{dv_z^+}{dy^+} = \frac{1}{ky^+} \Rightarrow v_z^+ = \frac{1}{k} \ln y^+ + C$$

$$\boxed{v_z^+ = \frac{1}{k} \ln y^+ + C \quad (y^+ \rightarrow \infty)}$$

* C is hard to get using this approximate solution method. to get it involves a more complicated asymptotic matching technique, or, finding the full solution without approximation. Deen gives

$$C = \frac{1}{k} [\ln(4k) - 1]$$

but this doesn't fit experiment well. Instead

$$v_z^+ = \frac{1}{k} \ln y^+ + C \text{ is used}$$

with $k \approx 0.4$ and $C \approx 5.5$ to fit experimental data.

Finally, let's put this back in more useful variables.

$$v_z^+ = \frac{1}{K} \ln y^+ + C$$

$$K = 0.4, \quad C = 5.5$$

$$v_z^+ = \frac{\langle v_z \rangle}{(T_0/\rho)^{1/2}} \quad y^+ = \frac{u_\tau}{\nu} y = \left(\frac{T_0}{\rho}\right)^{1/2} \frac{1}{\nu} (R-r)$$

$$\langle v_z \rangle = \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \frac{1}{K} \ln \left[\left(\frac{T_0}{\rho}\right)^{1/2} \frac{R-r}{\nu} \right] + C \right\}$$

$\nwarrow \quad \nu = \mu/\rho$

$$= \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \frac{1}{K} \ln \left[\frac{T_0^{1/2}}{\rho^{1/2}} \frac{\rho}{\mu} (R-r) \right] + C \right\}$$

$$\langle v_z \rangle = \frac{T_0^{1/2}}{\rho^{1/2}} \left\{ \frac{1}{K} \ln \left[\frac{T_0^{1/2} \rho^{1/2}}{\mu} (R-r) \right] + C \right\}$$

← in the notes
p. 23-5

$$K = 0.4, \quad C = 5.5$$

Normalized to the centerline value.

$$\langle v_z \rangle(r=0) = \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \frac{1}{K} \ln \left[\underbrace{\frac{T_0^{1/2} \rho^{1/2}}{\mu} R}_{R^+} \right] + C \right\}$$

$$R^+ = R / (\nu/u_\tau) = \frac{u_\tau}{\nu} R = \frac{T_0^{1/2}}{\rho^{1/2} \nu} = \frac{T_0^{1/2} \rho^{1/2}}{\mu} R$$

$$\langle v_z \rangle(r=0) = \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \frac{1}{K} \ln R^+ + C \right\}$$

note also: $\frac{T_0^{1/2} \rho^{1/2}}{\mu} (R-r) = \frac{T_0^{1/2} \rho^{1/2}}{\mu} R \left(1 - \frac{r}{R}\right) = R^+ \frac{y}{R}$

$$\frac{\langle v_z \rangle}{\langle v_z \rangle(r=0)} = \frac{\frac{1}{K} \ln \left[\frac{R^+ y}{R} \right] + C}{\frac{1}{K} \ln R^+ + C}$$

Eq. 10.5-19 in book, p. 275

* let's compare this to the HW:

$$u = \frac{2}{R^2} \int_0^R \langle v_z \rangle(r) r dr$$

$$= \frac{2}{R^2} \int_0^R \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \frac{1}{k} \ln \left[\frac{T_0^{1/2} \rho^{1/2}}{u} (R-r) \right] + C \right\} r dr$$

$$= \frac{2}{R^2} \left(\frac{T_0}{\rho}\right)^{1/2} \int_0^R \left\{ \frac{1}{k} \ln \left[R^+ \left(1 - \frac{r}{R}\right) \right] + C \right\} r dr$$

$$\text{let } \eta = 1 - r/R \quad \eta(0) = 1, \eta(R) = 0$$

$$d\eta = -\frac{1}{R} dr \quad r = R(1-\eta)$$

$$= \frac{2}{R^2} \left(\frac{T_0}{\rho}\right)^{1/2} \int_1^0 \left\{ \frac{1}{k} \ln [R^+ \eta] + C \right\} R(1-\eta) (-1) d\eta$$

$(-1) \Rightarrow$ flip
integral
bounds

$$= 2 \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \int_0^1 \frac{1}{k} \ln [R^+ \eta] (1-\eta) d\eta + \int_0^1 C(1-\eta) d\eta \right\}$$

$$= 2 \left(\frac{T_0}{\rho}\right)^{1/2} \left\{ \int_0^1 \left(\frac{1}{k} \ln R^+ + C \right) (1-\eta) d\eta + \int_0^1 \frac{1}{k} \ln \eta (1-\eta) d\eta \right\}$$

$$\int_0^1 (1-\eta) d\eta = \eta - \frac{\eta^2}{2} \Big|_0^1 = 1 - \frac{1}{2} = \frac{1}{2}$$

mathematical

$$\int_0^1 (1-\eta) \ln \eta d\eta = -\eta + \frac{\eta^2}{4} + \eta \ln \eta - \frac{1}{2} \eta^2 \ln \eta \Big|_0^1 = -\frac{3}{4}$$

$$u = 2 \left(\frac{T_0}{p} \right)^{1/2} \left\{ \left(\frac{1}{k} \ln R^+ + c \right)^{1/2} + \frac{1}{k} (-3/4) \right\}$$

$$u = \left(\frac{T_0}{p} \right)^{1/2} \left\{ \frac{1}{k} \ln R^+ + c - \frac{3}{2} \frac{1}{k} \right\}$$

$$u = \left(\frac{T_0}{p} \right)^{1/2} \left[\underbrace{\frac{1}{k} \ln R^+}_{2.5} + \underbrace{c - \frac{3}{2k}}_{1.75} \right]$$

← similar to $v_z(r=0)$, but not quite the same.
 c vs. $c - \frac{3}{2k}$

$$u = \left(\frac{T_0}{p} \right)^{1/2} [2.5 \ln R^+ + 1.75]$$

Now, finally, what is $\frac{\langle v_z \rangle}{u} = ?$

$$\frac{\langle v_z \rangle}{u} = \frac{T_0^{1/2}/p^{1/2} \left\{ \frac{1}{k} \ln [R^+ (1 - \gamma/k)] + c \right\}}{T_0^{1/2}/p^{1/2} \left\{ \frac{1}{k} \ln R^+ + c - \frac{3}{2k} \right\}}$$

$$\frac{\langle v_z \rangle}{u} = \frac{\frac{1}{k} \ln [R^+ (1 - \gamma/k)] + c}{\frac{1}{k} \ln R^+ + c - \frac{3}{2k}}$$

$$k = 0.4$$

$$c = 9.5$$

$$R^+ = \frac{T_0^{1/2} p^{1/2}}{\mu} R$$