

1

sinusoidal pressure gradient:

$$\frac{\partial P}{\partial z} = -a \sin \omega t$$

↑ ↙

amplitude frequency

units: $\frac{Pa}{m} = \frac{kg}{m^2 s^2}$ units: $1/s$

$$\rho \left[\frac{\partial v_z}{\partial t} + \cancel{v_r \frac{\partial v_z}{\partial r}} + \cancel{\frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta}} + \cancel{v_z \frac{\partial v_z}{\partial z}} \right] =$$

$$-\frac{\partial p}{\partial z} + \mu \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \cancel{\frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2}} + \cancel{\frac{\partial^2 v_z}{\partial z^2}} \right]$$

$$N_r = 0, \quad N_\theta = 0, \quad \frac{\partial v_z}{\partial z} = \frac{\partial^2 v_z}{\partial z^2} = \frac{\partial^2 v_z}{\partial \theta^2} = 0$$

$$\frac{\partial v_z}{\partial t} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \frac{\mu}{\rho} \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right)$$

$$\frac{\partial v_z}{\partial t} = \frac{a}{\rho} \sin(\omega t) + \frac{\mu}{\rho} \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right)$$

$$v_z(r=R) = 0$$

(no slip)

$$\frac{\partial v_z}{\partial r}(r=0) = 0$$

(Symmetry)

$$v_z(t=0) = 0$$

initial stagnant.

(b) Non-dimensionalize using

$$\eta = r/R, \quad \tau = \frac{vt}{R^2} = \frac{\mu t}{\rho R^2}$$

↑
viscous diffusion time scale: R^2/ν

$$\theta = \frac{v_z R}{\nu}$$

$$\frac{\nu}{R} \frac{\partial^2}{\partial \eta^2} \frac{\partial \theta}{\partial \tau} = \frac{a}{\rho} \sin\left(\omega \frac{R^2}{\nu} \tau\right) + \frac{\nu}{R^2} \frac{1}{\eta} \frac{\partial}{\partial \eta} \left[R \eta \cdot \frac{1}{R} \frac{\partial}{\partial \eta} \left(\frac{\partial v}{R} \right) \right]$$

$$\frac{\nu^2}{R^3} \frac{\partial \theta}{\partial \tau} = \frac{a}{\rho} \sin\left(\omega \frac{R^2}{\nu} \tau\right) + \frac{\nu^2}{R^3} \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial \theta}{\partial \eta} \right)$$

$$\frac{\partial \theta}{\partial \tau} = \frac{R^3 a}{\rho \nu^2} \sin\left(\frac{\omega R^2}{\nu} \tau\right) + \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial \theta}{\partial \eta} \right)$$

*up to 3 bonus points

$$\text{let } \boxed{\varepsilon = \frac{\omega R^2}{\nu}} = \frac{R^2 \omega / \nu}{\omega} = \frac{\text{viscous diffusion time}}{\text{pressure time const.}}$$

① Strouhal number, St

$$\text{let } \boxed{b = \frac{R^3 a}{\rho \nu^2}} = \frac{R^2}{\mu(\nu/\rho)} \cdot a = \frac{R^2}{\mu \mu} a = \frac{\text{pressure gradient}}{\text{viscous stress gradient}}$$

$$\boxed{\frac{\partial \theta}{\partial \tau} = b \sin(\varepsilon \tau) + \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial \theta}{\partial \eta} \right)}$$

$$\boxed{\varepsilon = \frac{\omega R^2}{\nu}} \quad \boxed{b = \frac{R^3 a}{\rho \nu^2}}$$

Boundary & initial conditions

$$\boxed{\theta(\eta=1) = 0}$$

(no slip)

$$\boxed{\frac{\partial \theta}{\partial \eta}(\eta=0) = 0}$$

(symmetry)

$$\boxed{\theta(\tau=0) = 0}$$

(stagnant initially)

(c) verify that $\Phi = f(\eta) b e^{i\epsilon\tau}$ for the PDE

$$\frac{\partial \Phi}{\partial \tau} = b e^{i\epsilon\tau} + \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial \Phi}{\partial \eta} \right)$$

and find an ODE for $f(\eta)$.

$$\frac{\partial}{\partial \tau} (f(\eta) b e^{i\epsilon\tau}) = b e^{i\epsilon\tau} + \frac{1}{\eta} \frac{\partial}{\partial \eta} \left[\eta \frac{\partial}{\partial \eta} (f(\eta) b e^{i\epsilon\tau}) \right]$$

$$i\epsilon \cancel{b e^{i\epsilon\tau}} f = \cancel{b e^{i\epsilon\tau}} + \cancel{b e^{i\epsilon\tau}} \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial f}{\partial \eta} \right)$$

* cancels, is a solution!

$$i\epsilon f = 1 + \frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial f}{\partial \eta} \right)$$

or

$$\boxed{\frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial f}{\partial \eta} \right) - i\epsilon f + 1 = 0}$$

Boundary conditions

$$\boxed{f(\eta=1)=0}$$

$$\boxed{\frac{\partial f}{\partial \eta}(\eta=0)=0}$$

probably more correct to use total derivatives now.

Either ok for Exam.

* Can also expand derivatives:

$$\frac{1}{\eta} \frac{\partial}{\partial \eta} \left(\eta \frac{\partial f}{\partial \eta} \right) = \frac{1}{\eta} \frac{\partial f}{\partial \eta} + \frac{\partial^2 f}{\partial \eta^2}$$

$$\boxed{\frac{\partial^2 f}{\partial \eta^2} + \frac{1}{\eta} \frac{\partial f}{\partial \eta} - i\epsilon f + 1 = 0}$$

(d) Solve Expression from (c) using perturbation theory.

$$f(\eta) = f_0 + \varepsilon f_1 + O(\varepsilon^2)$$

* Substitute into ODE:

$$\frac{d^2}{d\eta^2}(f_0 + \varepsilon f_1) + \frac{1}{\eta} \frac{d}{d\eta}(f_0 + \varepsilon f_1) - i\varepsilon(f_0 + \varepsilon f_1) + 1 = 0$$

at order 0: $\frac{d^2 f_0}{d\eta^2} + \frac{1}{\eta} \frac{df_0}{d\eta} + 1 = 0 \quad \leftarrow \text{no } i\varepsilon f_0$

at order 1: $\varepsilon \frac{d^2 f_1}{d\eta^2} + \varepsilon \frac{1}{\eta} \frac{df_1}{d\eta} - i\varepsilon f_0 = 0 \quad \leftarrow \text{no } 1$

* Boundary conditions:

$$f_0(\eta=1) + \varepsilon f_1(\eta=1) = 0 \Rightarrow f_0(\eta=1) = 0$$

$$f_1(\eta=1) = 0$$

$$\frac{df_0}{d\eta}(\eta=0) + \varepsilon \frac{df_1}{d\eta}(\eta=1) = 0$$

$$\Rightarrow \frac{df_0}{d\eta}(\eta=0) = 0$$

$$\frac{df_1}{d\eta}(\eta=0) = 0$$

* Solve at $O(\varepsilon^0)$:

$$\frac{d^2 f_0}{d\eta^2} + \frac{1}{\eta} \frac{df_0}{d\eta} = -1$$

$$\frac{1}{\eta} \frac{d}{d\eta} \left(\eta \frac{df_0}{d\eta} \right) = -1 \quad \rightarrow \text{integrate}$$

$$\eta \frac{df_0}{d\eta} = -\frac{\eta^2}{2} + c_1$$

$$\frac{df_0}{d\eta} = -\frac{\eta}{2} + \frac{c_1}{\eta} \quad \rightarrow \text{integrate}$$

$$f_0 = -\frac{\eta^2}{4} + C_1 \ln \eta + C_2$$

• apply BC's :

$$\frac{df_0}{d\eta}(\eta=0) \text{ diverges as } \eta \rightarrow 0, \text{ so } C_1 = 0$$

$$f_0(\eta=1) = -\frac{1}{4} + C_2 = 0 \Rightarrow C_2 = \frac{1}{4}$$

$$\boxed{f_0 = \frac{1}{4}(1-\eta^2)}$$

* Solve at $\theta(\tau')$:

$$\frac{d^2 f_1}{d\eta^2} + \frac{1}{\eta} \frac{df_1}{d\eta} - i f_0 = 0$$

$$\frac{1}{\eta} \frac{d}{d\eta} \left(\eta \frac{df_1}{d\eta} \right) = \frac{i}{4} (1-\eta^2)$$

$$\frac{d}{d\eta} \left(\eta \frac{df_1}{d\eta} \right) = \frac{i}{4} (\eta - \eta^3) \quad \downarrow \text{ integrate}$$

$$\eta \frac{df_1}{d\eta} = \frac{i}{4} \left(\frac{\eta^2}{2} - \frac{\eta^4}{4} \right) + C_1$$

$$\frac{df_1}{d\eta} = \frac{i}{4} \left(\frac{\eta}{2} - \frac{\eta^3}{4} \right) + \frac{C_1}{\eta} \quad \downarrow \text{ integrate}$$

$$f_1 = \frac{i}{4} \left(\frac{\eta^2}{4} - \frac{\eta^4}{16} \right) + C_1 \ln \eta + C_2$$

• Apply BC's

$$\frac{df_1}{d\eta}(\eta=0) \text{ diverges as } \eta \rightarrow 0, \text{ so } C_1 = 0$$

$$f_1(\eta=1) = \frac{i}{4} \left(\frac{1}{4} - \frac{1}{16} \right) + C_2 = 0$$

$$C_2 = -\frac{3i}{64}$$

$$\boxed{f_1 = \frac{i}{64} (4\eta^2 - \eta^4 - 3)}$$

* Put it all together:

$$f(\eta) = \frac{1}{4}(1-\eta^2) + \frac{i\varepsilon}{64}(4\eta^2 - \eta^4 - 3) + O(\varepsilon^2)$$

(e) Find an expression for $\theta(\eta, \tau)$:

$$\Phi = f(\eta) b e^{i\varepsilon\tau}$$

$$\Phi = \frac{b e^{i\varepsilon\tau}}{4}(1-\eta^2) + \frac{i\varepsilon b e^{i\varepsilon\tau}}{64}(4\eta^2 - \eta^4 - 3) + O(\varepsilon^2)$$

$$e^{i\varepsilon\tau} = \cos(\varepsilon\tau) + i \sin(\varepsilon\tau)$$

$$i\varepsilon e^{i\varepsilon\tau} = i\varepsilon \cos(\varepsilon\tau) + i^2 \varepsilon \sin(\varepsilon\tau)$$

$$= i\varepsilon \cos(\varepsilon\tau) - \varepsilon \sin(\varepsilon\tau)$$

$$\theta = \text{Im}(\Phi)$$

$$\theta = \frac{b \sin(\varepsilon\tau)}{4}(1-\eta^2) + \frac{b\varepsilon \cos(\varepsilon\tau)}{64}(4\eta^2 - \eta^4 - 3) + O(\varepsilon^2)$$

* The zero order term is in phase. ①

* up to 3 bars points

The first order term is out of phase ②

* The zero order term is traditional poiseuille ③

flow, but modulated by $b \sin(\varepsilon\tau)$.