

Chemical Engineering 374

Fluid Mechanics

NonNewtonian Fluids



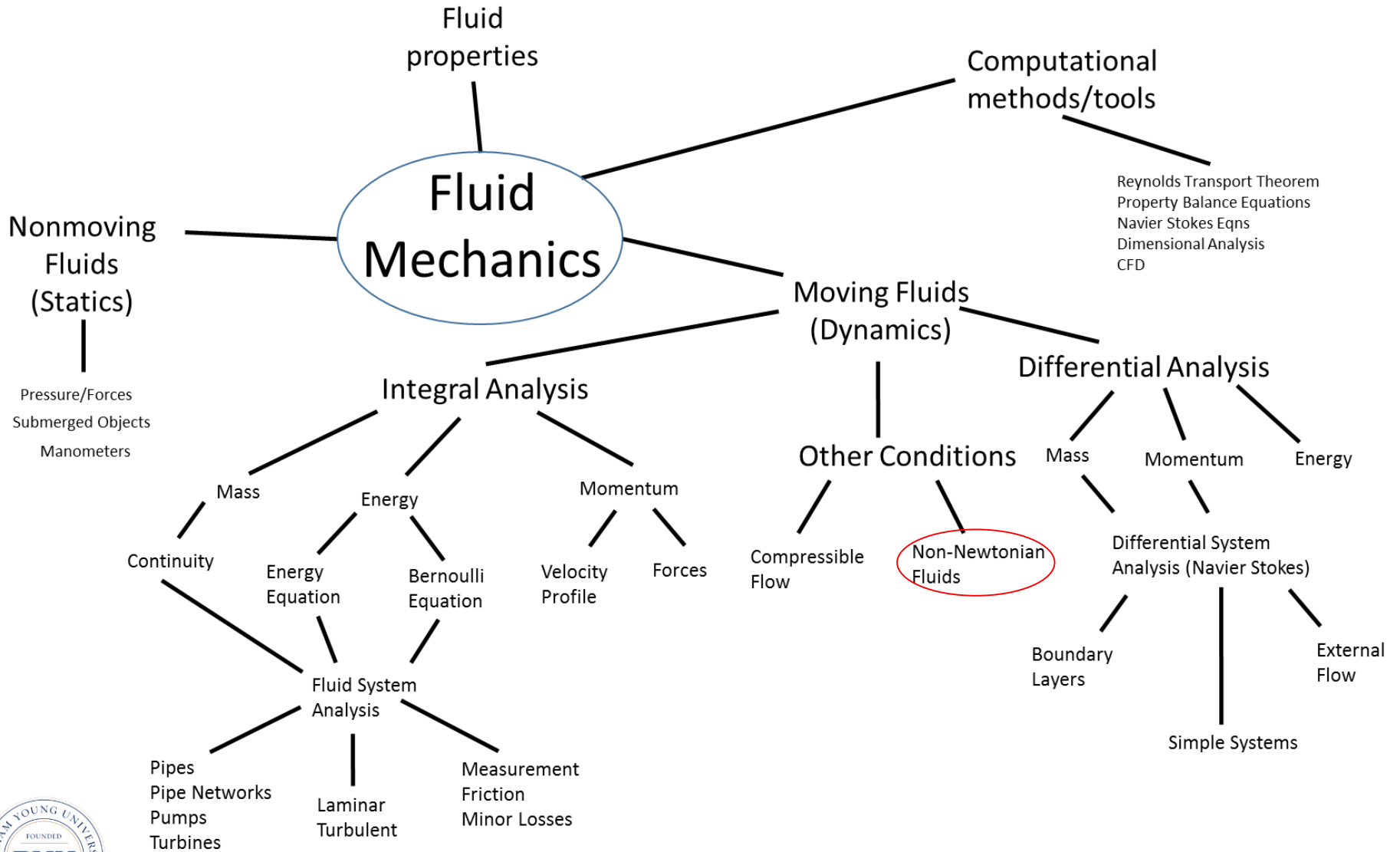
Spiritual Thought

“The recent emphasis of making the Sabbath a delight is a direct result of inspiration from the Lord.”

M. Russell Ballard



Fluids Roadmap



OEP 10 Clip



OEP 10

Open Ended Problem #10

The Avengers

INDIVIDUAL WORK ONLY, Due 11/30/16 at beginning of class
(Don't be afraid to "Google" good assumptions!)

The Avengers

In a thematic scene of pure awesome, the Avengers stop Loki and his army of Chitauri from conquering earth though a wormhole opened by the Tesseract (we won't go into the inaccuracies of the nuke and the blast in the Chitauri mother-ship... in this class) successfully closing the portal and saving the earth from alien invaders (how cliché). However, there is a serious problem with this scene. Although mass is supposed to be moving through the portal, no air is going through the wormhole, which it would be if a wormhole were opened in the atmosphere of earth! Given this, calculate the rate at which air passes through the portal when it is first opened.



Key Points

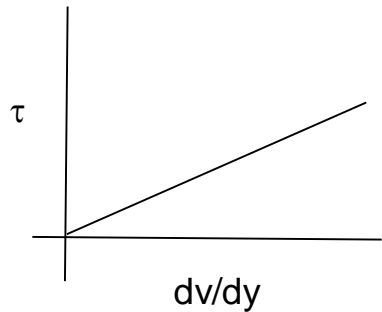
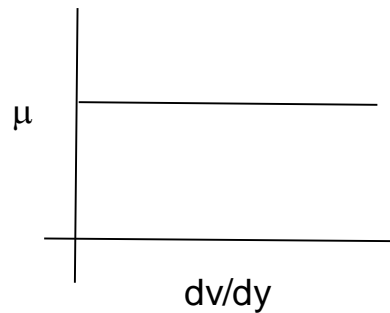
- Types and properties of non-Newtonian Fluids
- Pipe flows for non-Newtonian fluids
- Velocity profile / flow rate
- Pressure drop
 - Friction factor
 - Pump power
- Rheological Parameters

Power Law Fluids



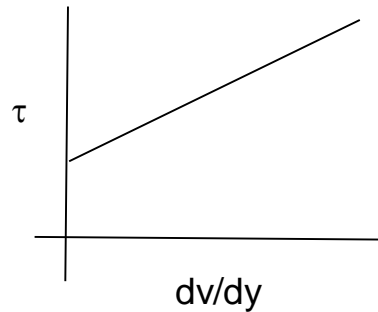
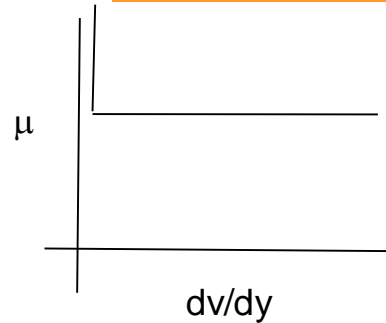
Non-Newtonian Fluids

Newtonian



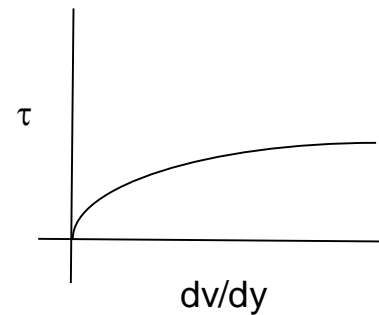
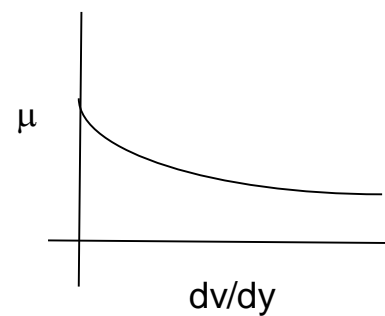
$$\tau = \mu * dv/dy$$

Bingham Plastic



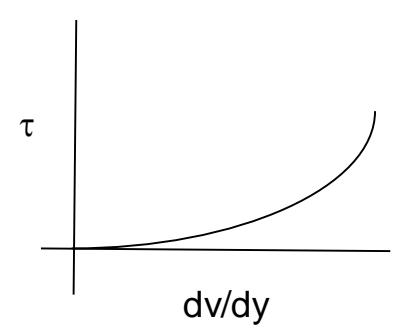
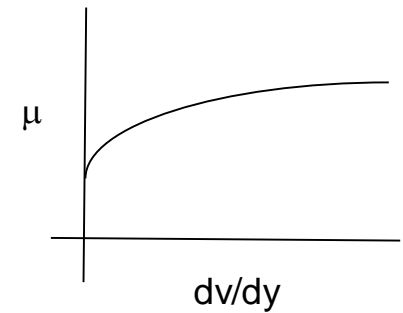
$$\tau = \mu * dv/dy + \tau_y$$

Pseudoplastic



$$\tau = \kappa * |dv/dy|^n$$

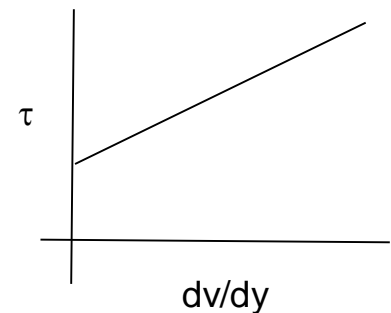
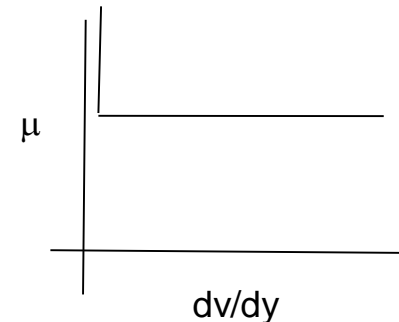
Dilatant



Bingham Plastic

- 3D elastic structures. Weak solid structures that must be broken
- Resists small shear, but structure “breaks apart” with large shear.
- Then τ is \sim linear with du/dx
- Some slurries (coal, grain slurries), sewage sludge.
 - Toothpaste (no drip)
- Larger particles \rightarrow weak solid structure \rightarrow breaks

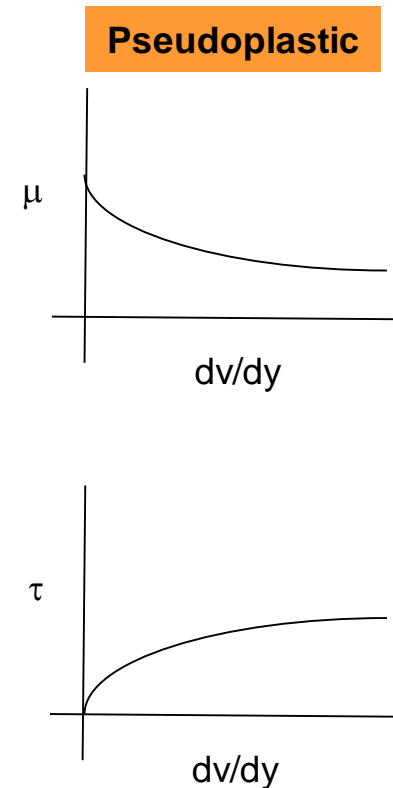
Bingham Plastic



$$\tau = \mu * dv/dy + \tau_y$$

Pseudoplastic

- Most common
- Dissolved or dispersed particles, like dissolved long chain molecules.
- Have a random orientation in the fluid at rest, but line up when the fluid is sheared.
 - τ decreases with strain rate
 - μ drops as molecules align
- Polymer melts, paper pulp suspensions, pigment suspensions, hair gel, blood, muds, most slurries
- “Shear-Thinning”
 - motor oil

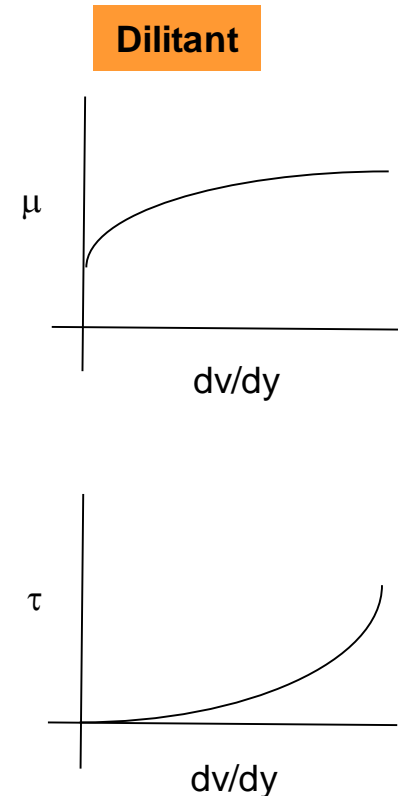


$$\tau = K * (-dv/dy)^n$$



Dilatant

- Rare
- Slurries of solid particles with barely enough liquid to keep apart. (corn starch, water squeezed out at high shear)
- At low strain rates, the fluid can lubricate solids; at high strain rates, this lubrication breaks down.
- μ increases with strain $\rightarrow \tau$ increases.
- “Shear thickening”.



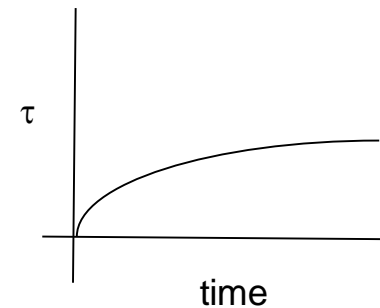
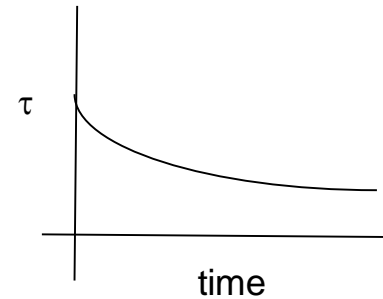
$$\tau = K * (-dv/dy)^n$$

Dilatant Example



Time dependence

- Thixotropic
 - Slurries/solutions of polymers
 - Many known fluids
 - Most are pseudoplastic
 - Alignable particles/molecules with weak bonds (H-bonding)
 - Paint
- Rheopectic
 - Rare
 - Fewer known examples
 - Usually fluids only show this behavior under mild shearing
- Changes occur within the first 60 sec. for most processes.
- Hard to describe
- Viscoelastic



Power Law Fluids

- Governing equations are “correct” in terms of τ
 - Expression for τ is the model.
 - Called a “constitutive relation”
 - Also have these for mass and heat fluxes in heat and mass transfer.

- Newtonian flow

$$\tau = -\mu \frac{dv}{dy}$$

- For dilatant and pseudoplastic fluids (most common)—**Power Law**

$$\tau = K \left(-\frac{dv}{dy} \right)^n$$

$n > 1 \rightarrow$ Dilatant

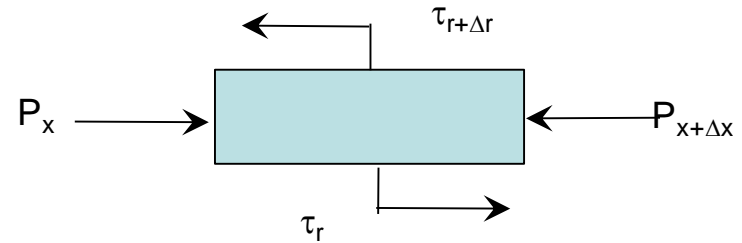
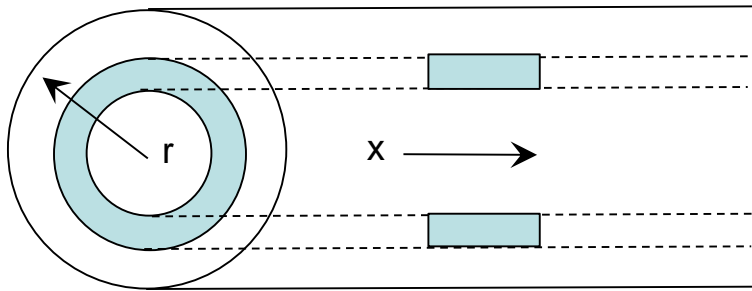
$n < 1 \rightarrow$ Pseudoplastic

$n = 1, K = \mu \rightarrow$ Newtonian

- K, n are empirical constants
- Many other forms
 - Simpler ones have 3 parameters and give a better fit, but are more complex than power law form.
 - See Handout of Book Chapter on Webpage.



Laminar Pipe Flow



Force Balance: Pressure, stress

$$(P_x - P_{x+\Delta x})(2\pi r \Delta r) + (2\pi \Delta x r) \tau_r - (2\pi \Delta x)(r + \Delta r) \tau_{r+\Delta r} = 0$$

Divide $2\pi \Delta r \Delta x$

$$r \frac{P_x - P_{x+\Delta x}}{\Delta x} + \frac{r \tau_r - (r + \Delta r) \tau_{r+\Delta r}}{\Delta r} = 0$$

Limit $\Delta x, \Delta r \rightarrow 0$

$$-\frac{dP}{dx} = \frac{1}{r} \frac{d(r\tau)}{dr} = C$$

Separate variables and integrate with $\tau=0$ at $r=0$

$$\tau = -\frac{r}{2} \frac{dP}{dx}$$

Non-Newtonian Pipe Flow

- Most non-Newtonian flows are laminar.
- Key results: (remember, Q is just volumetric flow rate- \dot{V})

- Force balance:
$$\tau = -\frac{r}{2} \frac{dP}{dx}$$
- Power law constitutive relation
$$\tau = K \left(-\frac{dv}{dr} \right)^n$$
- Integrate with B.C. $v=0$ at $r=R$
$$v = \left(-\frac{1}{2K} \frac{dP}{dx} \right)^{1/n} \left(\frac{n}{n+1} \right) \left(R^{\frac{n+1}{n}} - r^{\frac{n+1}{n}} \right)$$

- $\dot{V} = Q = A v_{avg}$
$$Q = \frac{\pi n D^3}{8(3n+1)} \left(-\frac{D}{4K} \frac{dP}{dx} \right)^{1/n} \quad V_{avg} = \frac{nD}{2(3n+1)} \left(-\frac{D}{4K} \frac{dP}{dx} \right)^{1/n}$$

- Kinetic Energy Correction Factor:
$$\alpha = \frac{3(3n+1)^2}{(5n+3)(2n+1)}$$

- Momentum Flux Correction Factor:
$$\beta = \frac{3n+1}{2n+1}$$



Pressure Drop—Laminar Flow

Define $f = \frac{4\tau_w}{\frac{1}{2}\rho v_{avg}^2}$ 1

Force Balance $\frac{dP}{dx} = -\frac{2\tau_w}{R}$ 2

$$V_{avg} = -\frac{1}{\pi R^2} \int_A v(r) dA$$

Newtonian

$$v = -\frac{R^2}{4\mu} \left(\frac{dP}{dx} \right) \left(1 - \frac{r^2}{R^2} \right)$$

3

$$V_{avg} = -\frac{R^2}{8\mu} \left(\frac{dP}{dx} \right)$$

4

$$\tau_w = \frac{4\mu V_{avg}}{R}$$

$$f = \frac{64}{Re}$$

Solve 4 for dP/dx
and insert into 2

Non-Newtonian

$$v = \left(-\frac{1}{2K} \frac{dP}{dx} \right)^{1/n} \left(\frac{n}{n+1} \right) \left(R^{\frac{n+1}{n}} - r^{\frac{n+1}{n}} \right)$$

$$V_{avg} = \frac{nD}{2(3n+1)} \left(-\frac{D}{4K} \frac{dP}{dx} \right)^{1/n}$$

$\tau_w =$ (some complex expression)

Insert into 1 for τ_w

$$f = \frac{8K}{\rho V_{avg}^2} \left(\frac{2(3n+1)V_{avg}}{nD} \right)^n$$



Turbulent Flow

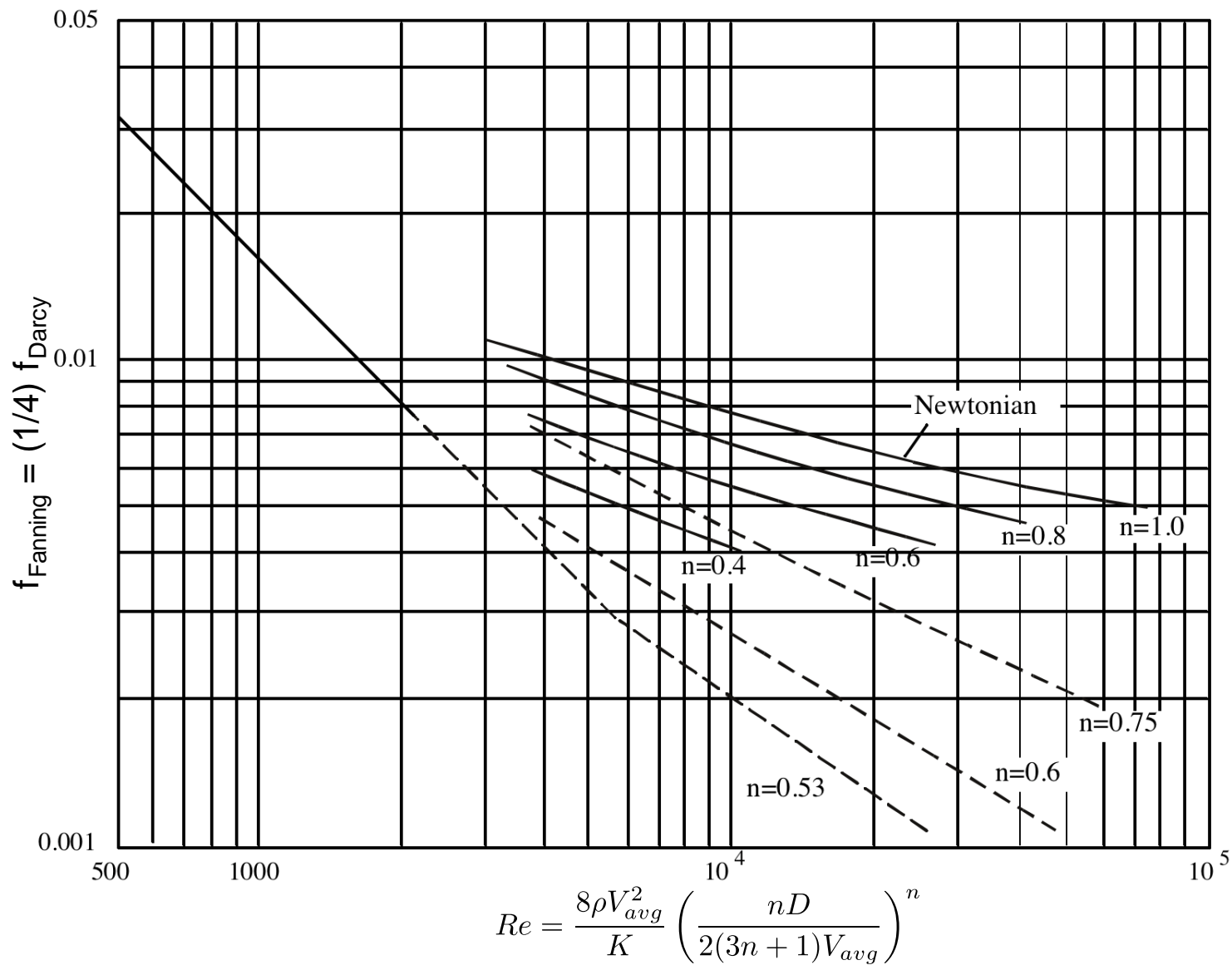
- Define the friction factor as before: $f = \frac{8\tau_w}{\rho V_{avg}^2} = \frac{\Delta PD/L}{\frac{1}{2}\rho V_{avg}^2}$
 - (Laminar or Turbulent)
- For turbulent flow we had $f = f(Re, \varepsilon/D)$ from dimensional analysis.
- Question: Will this work for non-Newtonian Flow?
- Question: What is the Reynolds number?
 - No clear definition of Re since μ is not constant (depends on the strain rate dv/dr , which depends on V_{avg})
- Use the same definition as the laminar friction factor: $Re=64/f$

$$f = \frac{8K}{\rho V_{avg}^2} \left(\frac{2(3n+1)V_{avg}}{nD} \right)^n \longrightarrow Re = \frac{8\rho V_{avg}^2}{K} \left(\frac{nD}{2(3n+1)V_{avg}} \right)^n$$

- (Definition based on laminar Newtonian, but used for all regimes)
- Plot friction factor versus Re as for Newtonian flows, using the red definition of Re.



Non-Newtonian Friction Factor (Power Law)



Rheological Parameters (power law)

- **Problem:** Non-Newtonian fluid has:
 - How to find **K**, and **n** for a given fluid?
- You need to measure something (what?)
- Try a pipe flow
 - $D, Q, dP/dx$
- Here's what we know:

$$\tau = K \left(-\frac{dv}{dr} \right)^n$$

$$\tau = K \left(-\frac{dv}{dr} \right)^n$$

$$v = \left(-\frac{1}{2K} \frac{dP}{dx} \right)^{1/n} \left(\frac{n}{n+1} \right) \left(R^{\frac{n+1}{n}} - r^{\frac{n+1}{n}} \right)$$

$$\tau = -\frac{r}{2} \frac{dP}{dx}$$

$$Q = \frac{\pi n D^3}{8(3n+1)} \left(-\frac{D}{4K} \frac{dP}{dx} \right)^{1/n}$$

$$\tau_w = K \left(-\frac{dv}{dr} \right)_w^n$$

$$V_{avg} = \frac{nD}{2(3n+1)} \left(-\frac{D}{4K} \frac{dP}{dx} \right)^{1/n}$$

$$\tau_w = -\frac{R}{2} \frac{dP}{dx}$$

- $D, Q, dP/dx \rightarrow V_{avg}, \tau_w$.
- Then relate these to **K**, **n**:
 - Compute $(-dv/dr)_w$ from $v(r)$

$$\tau_w = K \left(-\frac{dv}{dr} \right)_w^n$$



Rheological Parameters (power law)

- From $v(r)$, we get: $\left(-\frac{dv}{dr}\right)_w = \frac{2(3n+1)V_{avg}}{nD}$
- Now $\tau = K \left(-\frac{dv}{dr}\right)^n \longrightarrow \ln(\tau_w) = \ln(K) + n \ln\left(-\frac{dv}{dr}\right)_w$
- So a plot of $\ln(\tau_w)$ versus $\ln\left(-\frac{dv}{dr}\right)_w$ is linear with slope n , and intercept $\ln(K)$.
- But, note that $\left(-\frac{dv}{dr}\right)_w$ involves n , which is unknown \rightarrow what to do?
- Just rearrange:

$$\ln(\tau_w) = \ln(K) + n \ln(2(3n+1)V_{avg}/nD)$$

$$\ln(\tau_w) = n \ln(V_{avg}) + \{\ln(K) + n \ln(2(3n+1)/nD)\}$$

- Now, a plot of $\ln(\tau_w)$ versus $\ln(V_{avg})$ is linear with slope n .
- Once n is known, K can be computed from the intercept (term in $\{\}$), or just compute it analytically from $\tau = K \left(-\frac{dv}{dr}\right)^n$ and $\left(-\frac{dv}{dr}\right)_w = \frac{2(3n+1)V_{avg}}{nD}$ which give

$$K = \frac{\tau_w}{(2(3n+1)V_{avg}/nD)^n}$$



Recap

- To compute K , n for a non-Newtonian fluid
- Measure Q , D , dP/dx
- Compute V_{avg} from Q and D (area), that is, $Q=A*V_{avg}$
- Compute τ_w from $\tau_w = -\frac{R}{2} \frac{dP}{dx}$
- Plot $\ln(\tau_w)$ versus $\ln(V_{avg})$
- Fit a line to the data (the linear part of the data)
- The slope is n
- K is computed from the intercept, or from

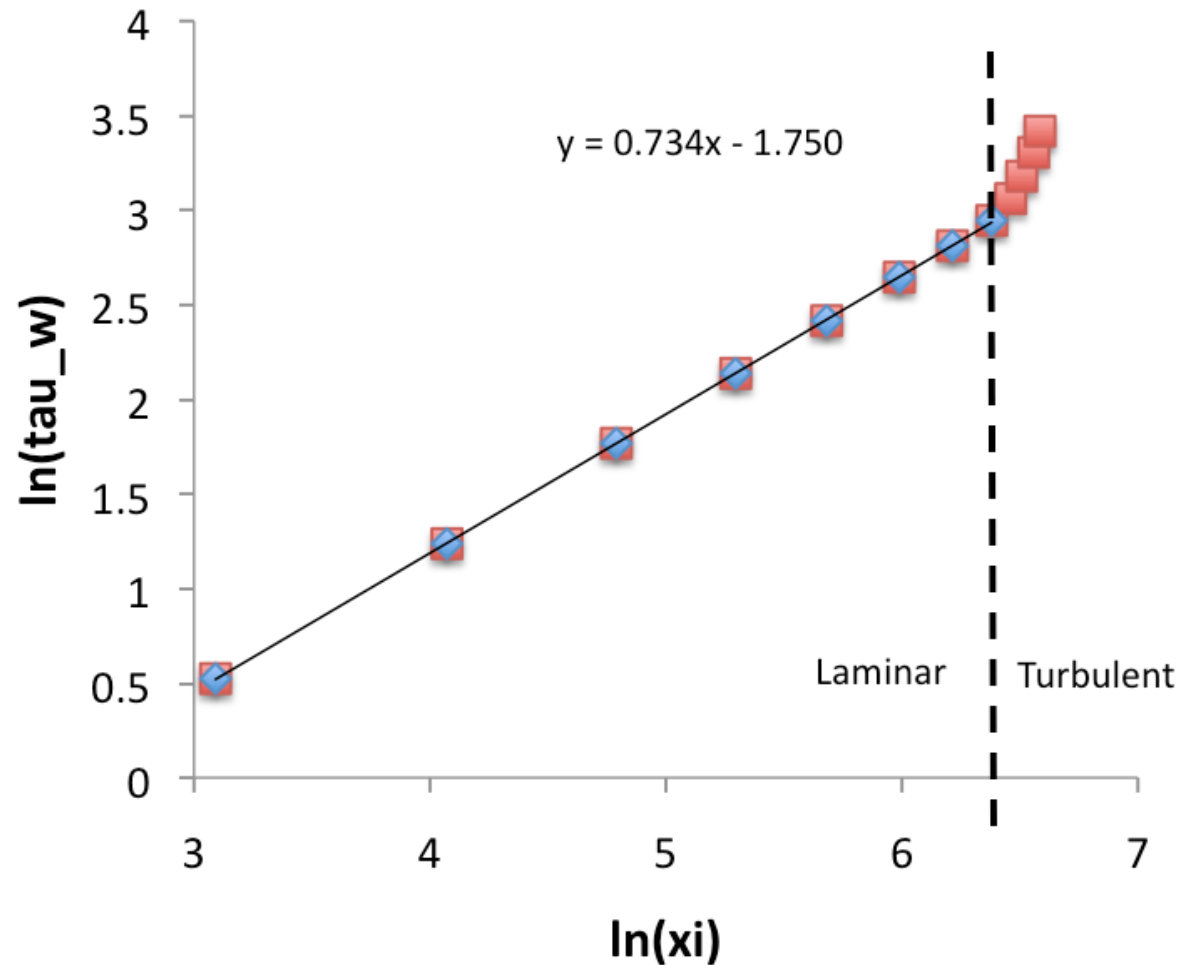
$$K = \frac{\tau_w}{(2(3n + 1)V_{avg}/nD)^n}$$

Note, the units on K are $(\text{kg} \cdot \text{s}^{n-2}/\text{m})$



Example

- Given:
 - Diameter
 - Pressure Drop
 - Flow Rate
- Compute:
 - K, n
 - Re
 - Power through a given pipe is as usual, $Q \cdot \Delta P$



Note, here $\xi = 8 \cdot V_{avg} / D$, and instead of plotting $\ln(\tau_w)$ versus $\ln(V_{avg})$, I'm plotting $\ln(\tau_w)$ versus $\ln(\xi)$. The approach is the same, but the intercept has a different formula for getting K. By the way, $\xi = 8 \cdot V_{avg} / D$ is $(-dv/dr)_w$ for Newtonian fluids, hence that choice here.

