Chemical Engineering 412

Introductory Nuclear Engineering

Lecture 7 Nuclear Fission Radiation Interactions with Matter



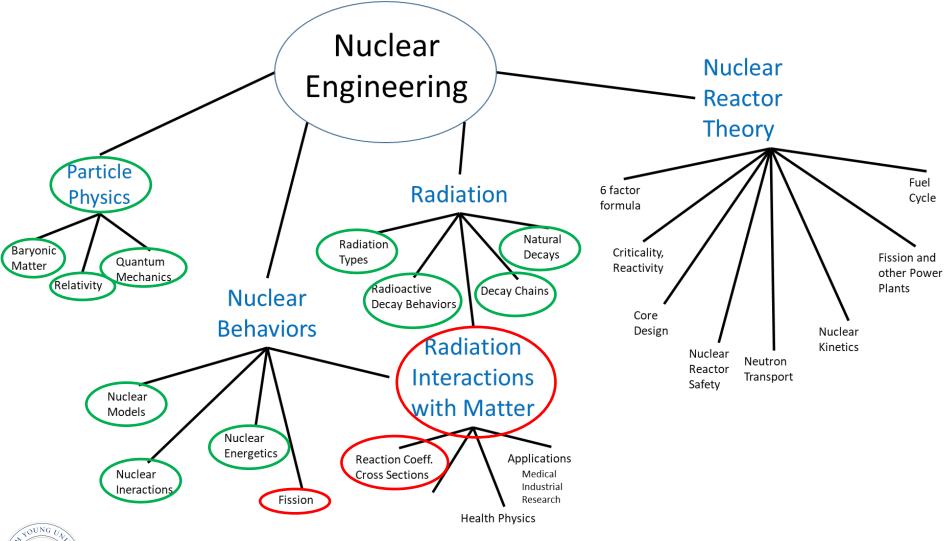
Spiritual Thought

"No man, having put his hand to the plough, and looking back, is fit for the kingdom of God." When difficult things are asked of us, even things contrary to the longings of our heart, remember that the loyalty we pledge to the cause of Christ is to be the supreme devotion of our lives. Although Isaiah reassures us it is available "without money and without price"—and it is—we must be prepared, using T. S. Eliot's line, to have it cost "not less than everything."



Jeffery R. Holland

Roadmap





- Understand mechanics and details of Fission
- Understand how radiation interacts with Matter
- Be able to calculate probabilities of interaction and radiation field intensities
- Understand both linear interaction coefficients and cross-sections
- Be able to calculate or find $\mu,\,\sigma,\,and\,\Sigma$



Neutron Interactions

- Elastic scattering (n,n) collision with no reaction and no change in total kinetic energies. Energy neutral.
- Inelastic scattering (n,n') collisions with energy absorption by nucleus. endoergic
- Radiative capture (n, γ) Capture of neutron by nucleus followed by γ-ray emission. excergic.
- Charged particle reactions (n,α) Neutron reaction to form α particles or protons. endoergic and exoergic.
- Neutron producing reactions (n,xn) Reactions with a net increase in neutrons. endoergic. (n,2n) important for ²H and ⁹Be.
- Fission (n,) forms multiple products Nucleus forms daughters. Generally exoergic.



Fission Reactions

 $^{235}_{92}U$ is <u>fissile</u> (undergoes fission)

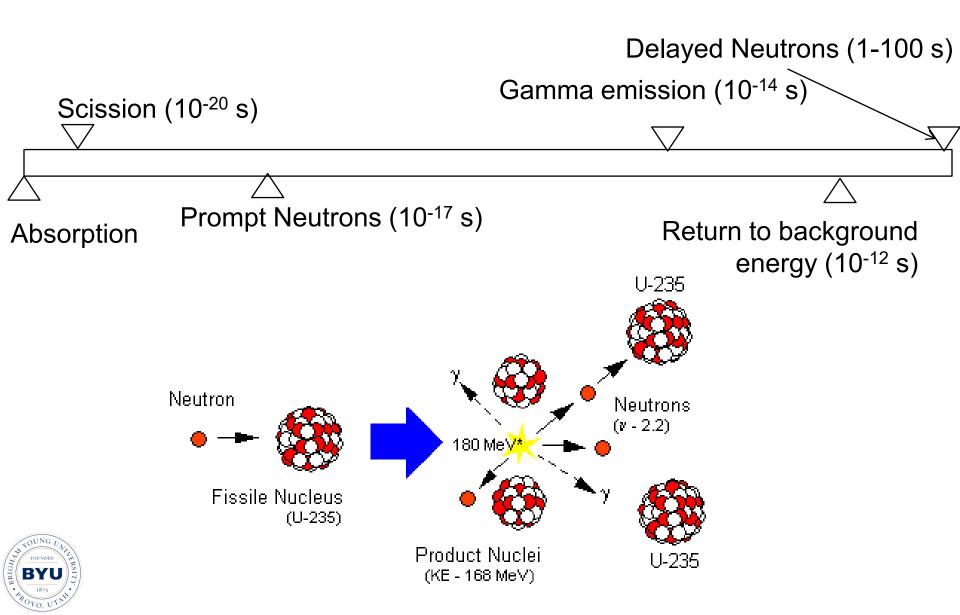
 $^{238}_{92}U$ is <u>fertile</u> (converts to a fissionable isotope)

Possible outcomes of $^{235}_{92}U$ reaction with neutron

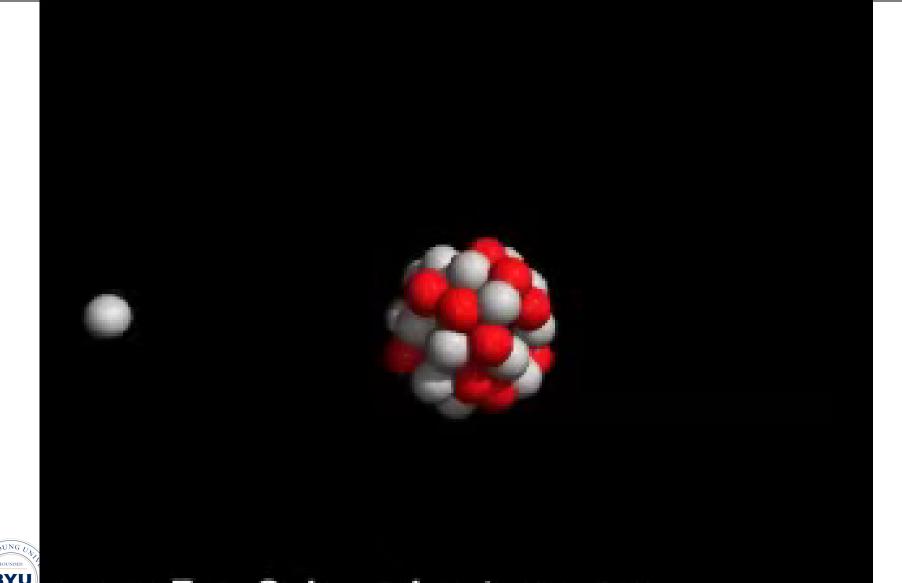
$$n + {}^{235}_{92}U \rightarrow \begin{cases} {}^{235}_{92}U + n & elastic \ scatter \\ {}^{235}_{92}U + n + \gamma & inelastic \ scatter \\ {}^{236}_{92}U + \gamma & radiative \ capture \\ {}^{236}_{92}U + \gamma & radiative \ capture \\ {}^{Y}_{H} + {}^{Y}_{L} + {}^{Y}_{1} + {}^{Y}_{2} + \cdots & fission \end{cases}$$



Fission (logarithmic) Timeline



Fission Animation





www.FreeScienceLectures.com





Emitted/Recoverable Energy

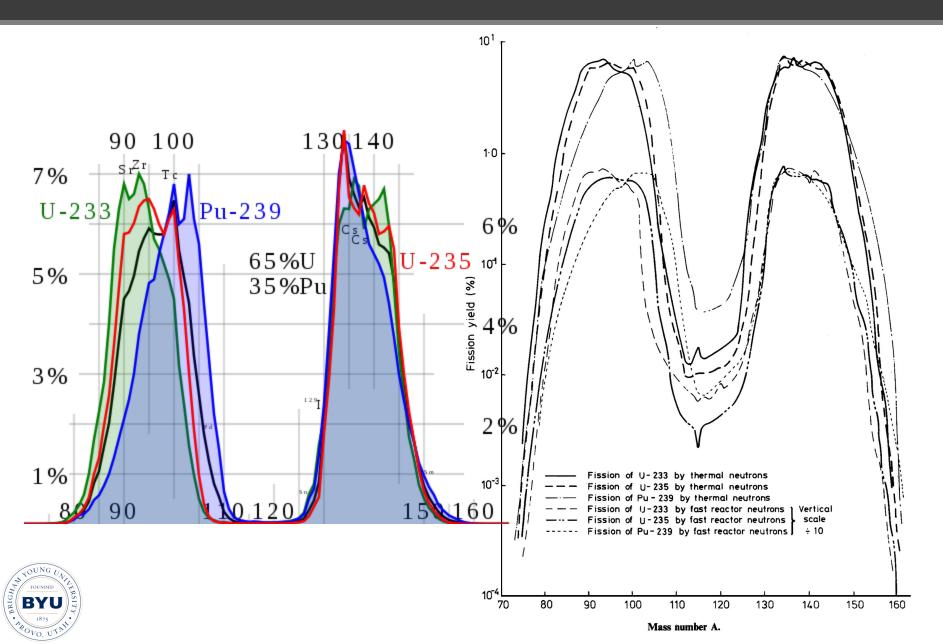
TABLE 3.6 EMITTED AND RECOVERABLE ENERGIES FOR FISSION OF ²³⁵U

Form	Emitted Energy, Me V	Recoverable Energy, Me V
Fission fragments	168	168
Fission-product decay		
β -rays	8	8
γ-rays	7	7
neutrinos	12	
Prompt γ -rays	7	7
Fission neutrons (kinetic energy)	5	5
Capture γ -rays	—	3-12
Total	207	198–207

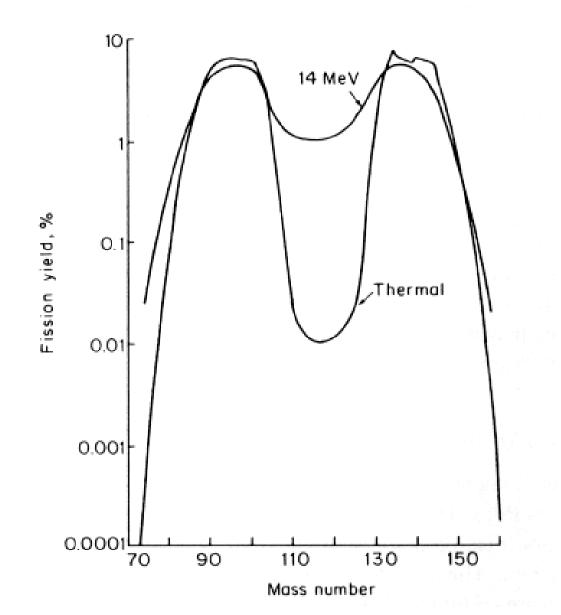


 $E_f = 200 \text{ MeV}$

Fission Product Distribution



Product Distribution at High Energy





Delayed Neutrons

- A small fraction (<1%) of total neutron production occur seconds or minutes after scission, represented by β below. These delayed neutrons are essential to reactor control.
- Fast neutron emission alone is far too rapid to allow control.
 Fast Fission Thermal Fission

Fast Fission		Thermal Fission		
$\overline{ u}$	β	$\overline{\overline{\nu}}$	β	
2.57	0.0064	2.43	0.0065	
2.62	0.0026	2.48	0.0026	
3.09	0.0020	2.87	0.0021	
	_	3.14	0.0049	
2.79	0.0148			
2.44	0.0203		1 (
3.3	0.0026	-		
	$\overline{ u}$ 2.57 2.62 3.09 - 2.79 2.44	$\begin{array}{c c} \overline{\nu} & \beta \\ \hline 2.57 & 0.0064 \\ 2.62 & 0.0026 \\ \hline 3.09 & 0.0020 \\ \hline - & - \\ 2.79 & 0.0148 \\ 2.44 & 0.0203 \end{array}$	$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	

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Source: Keepin [1965].

Delayed Neutron Data

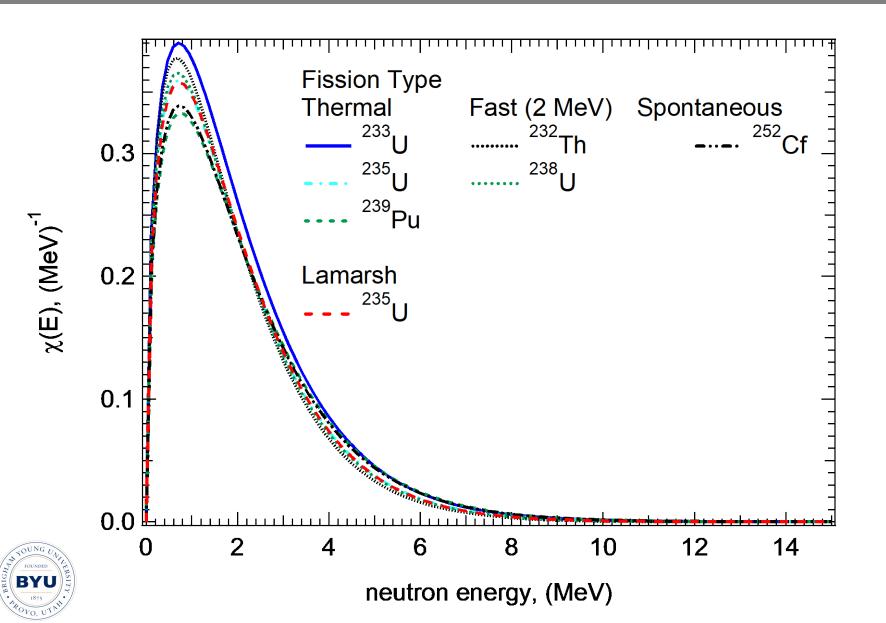
TABLE 3.5 DELAYED NEUTRON DATA FOR THERMAL FISSION IN ²³⁵U*

Group	Half-Life (sec)	Decay Constant (l_i, \sec^{-1})	Energy (ke V)	Yield, Neutrons per Fission	Fraction (β_i)
1	55.72	0.0124	250	0.00052	0.000215
2	22.72	0.0305	560	0.00346	0.001424
3	6.22	0.111	405	0.00310	0.001274
4	2.30	0.301	450	0.00624	0.002568
5	0.610	1.14		0.00182	0.000748
6	0.230	3.01		0.00066	0.000273
				Total	yield: 0.0158
				Total delayed fractio	

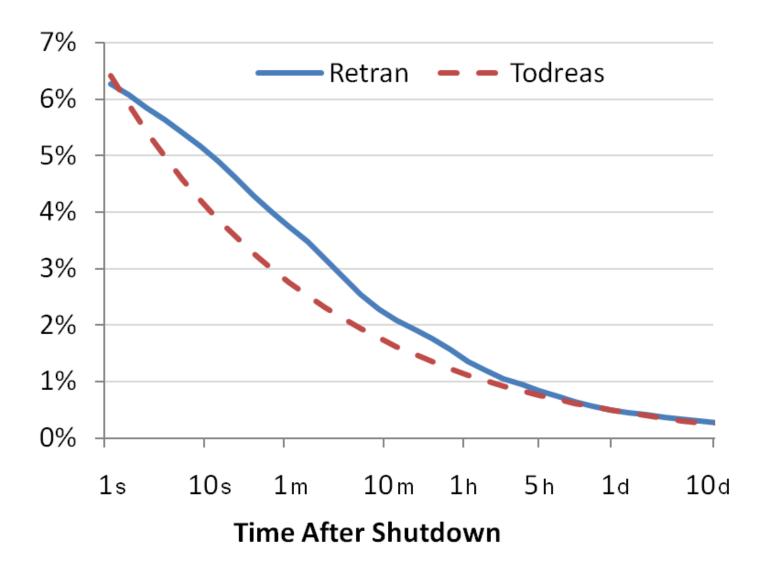
*Based in part on G. R. Keepin, *Physics of Nuclear Kinetics*, Reading, Mass.: Addison-Wesley, 1965.



Neutron Energy Spectrum



Decay Heat

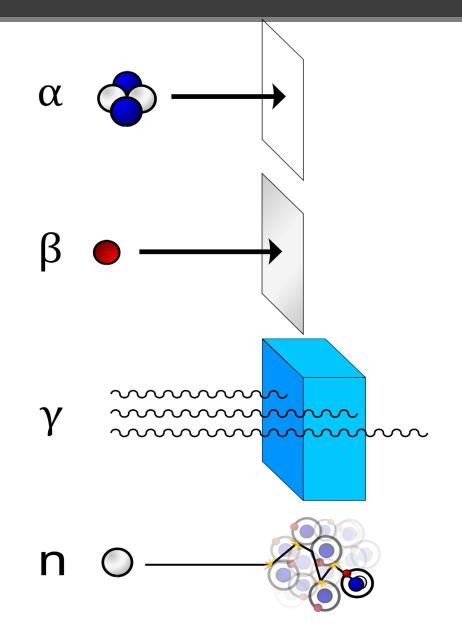




Radiation Interaction with Matter

- World is awash with radiation
- First step to understanding impact is knowing how it interacts
- Different particles have different effects
- Derive general terms to quantify interactions





Linear Interaction Coefficient

- As a particle passes through a homogeneous material
 - Probability of interaction is constant per differential unit distance traveled
 - Empirically derived

$$\mu_i \equiv \lim_{\Delta x \to 0} \frac{P_i(\Delta x)}{\Delta x}$$

- μ_i is called the macroscopic interaction coefficient
- indicated by Σ_i (except for photons).
- Depends on
 - Particle energy
 - Reaction Type
 - Scattering, absorption, fission, etc.
 - energy-dependent macroscopic linear absorption coefficient
 - linear fission coefficient
 - linear scattering coefficient, etc.
 - Medium type



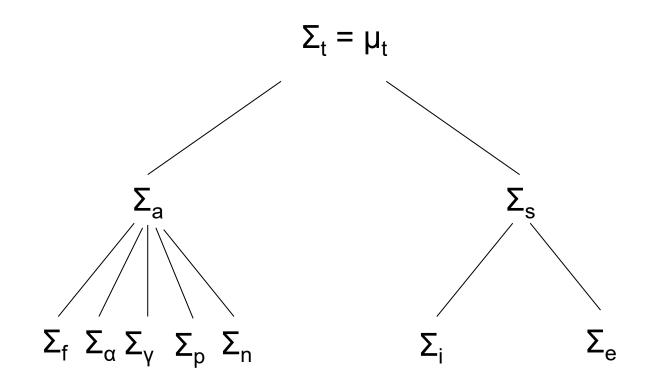
Total Probability of Interaction

- Interaction coefficients are divided into subcategories
 - i.e. total scattering coefficient, Σ_t
 - Linear scattering coefficients, Σ_s
 - Non-linear scattering coefficients
 - Total absorption coefficient, Σ_a
 - Neutron capture
 - Fission
 - Other absorbing interactions
- Total is sum of components
 - Radiation linear attenuation coefficient
 - Neutrons Cross Section
 - Photons Mass Interaction Coefficient

$$\mu_t(E) = \sum_i \mu_i(E)$$

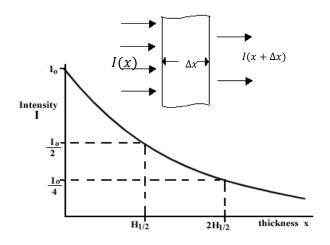


Or in other words...





Interaction in Material



The fractional amount of a beam that interacts in a differential slice of a material is given by

$$\frac{I(x) - I(x + \Delta x)}{I(x)} = P(x)$$

$$\mu_t \equiv \lim_{\Delta x \to 0} \frac{P_i(\Delta x)}{\Delta x} = \lim_{\Delta x \to 0} \frac{I(x) - I(x + \Delta x)}{\Delta x I(x)} = -\frac{1}{I(x)} \frac{dI(x)}{dx}$$

$$\frac{dI(x)}{dx} = -\mu_i I(x) \Rightarrow I(x) = I(0) \exp(-\mu_i x)$$



Conceptual Interpretations

- The linear attenuation coefficient can be thought of in three ways:
 - Probability that a particle interacts in a differential length of material (does not assume constant μ_i)
 - Inverse of the mean free path of a particle (assumes constant μ_i).
 - Related to distance at which half of particles have interacted $(x_{1/2, i} = \frac{\ln 2}{\mu_i})$ (assumes constant μ_i)
- Analogous to decay constants
 - Decay probabilities
 - Average lifetimes



- Half lives.

Non-absorbing Particles

- In many cases (scattering, photons, etc.), interactions do not eliminate the particles
- The total amount of particles
 - Highly complex, calculated with large computations
 - Derive a buildup factor, B(x), that correlates complex behavior with simple expression

 $I(x) = B(x)I(0)\exp(-\mu_i x)$

 This is especially common in calculating dose (as opposed to total particles).



Microscopic Cross Section

• Probability of interaction is proportional to the concentration of interaction sites/atoms

$$\mu_i = \sum_i = N\sigma_i = \sigma_i \frac{\rho N_a}{A}$$

- σ_i = microscopic cross section, has units of L²
- *N* = Number/atom density
- ρ = Mass density
- N_a = Avagadro's number
- *A* = Atomic mass of the medium



Cross sections for each interaction

$$\sigma_t = \sigma_e + \sigma_i + \sigma_\gamma + \sigma_f + \dots$$

total cross section

 $\sigma_{a} = \sigma_{\gamma} + \sigma_{f} + \sigma_{\alpha} + \sigma_{p} + \dots$

absorption cross section

scattering cross section

 $\sigma_t = \sigma_s + \sigma_a$

 $\sigma_s = \sigma_e + \sigma_i$

total cross section

t = total e = elastic scattering i = inelastic scattering γ = radiative capture f = fission α = alpha (charged) particle p = proton (charged) particle



Microscopic cross section

- The microscopic cross section
 - Independent of atomic density
 - Based strongly and complexly on particle kinetic energy
 - Play vital roles in nuclear engineering
- Behaviors are empirical!
 - (can be conceptually explained but not always quantitatively predicted by theoretical means)
- Typical unit is barns (1 barn = $1 \times 10^{-24} \text{ cm}^2$)
- 1 barn is approximate physical cross section of a uranium nucleus.

